

# Enhancing the deep-red/near-infrared fluorescence of higher rylene diimides *via* the chalcogen-annulation strategy

Kai Chen<sup>1</sup>, Xiao Chen<sup>2</sup>, Ke Hu<sup>3</sup>, Yilun Zhao<sup>1</sup>, Yujian Liu<sup>2</sup>, Guogang Liu<sup>1\*</sup>, Jinquan Chen<sup>3</sup>, Wei Jiang<sup>1,2</sup>, Zhigang Shuai<sup>2</sup>, Da-Hui Qu<sup>1</sup> & Zhaohui Wang<sup>1,2\*</sup>

<sup>1</sup>Key Laboratory for Advanced Materials and Joint International Research Laboratory of Precision Chemistry and Molecular Engineering, Feringa Nobel Prize Scientist Joint Research Center, Frontiers Science Center for Materiobiology and Dynamic Chemistry, Institute of Fine Chemicals, School of Chemistry and Molecular Engineering, East China University of Science and Technology, Shanghai 200237, China;

<sup>2</sup>Key Laboratory of Organic Optoelectronics and Molecular Engineering, Department of Chemistry, Tsinghua University, Beijing 100084, China;

<sup>3</sup>State Key Laboratory of Precision Spectroscopy, East China Normal University, 500 Dongchuan Road, Shanghai 200241, China

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Deep-red/near-infrared fluorescence is highly suitable for bioimaging owing to its ability to deeply penetrate tissues, organs, and live animals. However, developing organic fluorophores with high deep-red/near-infrared fluorescence quantum yield ( $\Phi_{\text{FL}}$ ) and fluorescent brightness remain a significant challenge owing to the energy gap law. Herein, we developed a straightforward and effective chalcogen-annulation strategy by introducing O, S and Se into the bay region of **TDI** and **QDI** fluorophores, realizing the increase of  $\Phi_{\text{FL}}$  and fluorescent brightness up to 10 times. To our best knowledge, this study potentially stands as the pioneering instance showcasing the anti-heavy-atom effect of chalcogens, and the absolute  $\Phi_{\text{FL}}$  (93%) and fluorescent brightness ( $128,200 \text{ cm}^{-1} \text{ mol}^{-1} \text{ L}$ ) of **Se-TDI** is among top deep-red/near-infrared organic fluorophores currently available. The femto-second transient absorption (fs-TA) measurements show the absence of obvious changes of the excited state lifetime after the introduction of chalcogens in **TDI** and **QDI** fluorophores, indicating that intersystem crossing (ISC) can be neglected in TDI and QDI fluorophores. Theoretical calculations further reveal the chalcogen-annulation strategy increase the radiative rates and reduce the reorganization energy of several accepting modes at the ground state in **TDI** fluorophores, leading to the suppression of internal conversion (IC) processes. Our chalcogen-annulation strategy, which effectively increases the  $\Phi_{\text{FL}}$  and restricts the IC processes, while remaining unaffected by the heavy-atom effect, offers novel insights and theoretical support for the design and synthesis of deep-red/near-infrared organic fluorophores with high  $\Phi_{\text{FL}}$  and fluorescent brightness.

**$\pi$ -extended rylene diimides, chalcogen-annulation strategy, deep-red/near-infrared fluorescence**

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## 1 Introduction

In recent years, there has been an increasing emphasis on deep-red/near-infrared fluorescence in the field of bioimaging, owing to its exceptional capacity to penetrate orga-

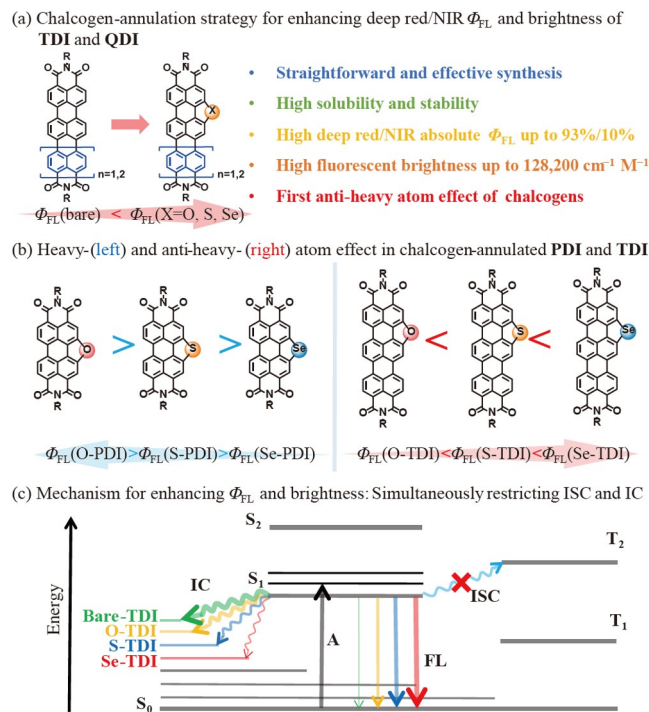
noids, tissues, organs, and live animals deeply while maintaining a high signal-to-background ratio [1–4]. Among the various deep-red/near-infrared materials, organic fluorophores are highly favored for their excellent derivatization, versatility, and well-established safety profile accumulated through years of clinical use [5–24]. However, developing deep-red/near-infrared organic fluorophores with high fluorescence quantum yield ( $\Phi_{\text{FL}}$ ) and fluorescent brightness

\*Corresponding authors (email: [liuguogang@ecust.edu.cn](mailto:liuguogang@ecust.edu.cn); [wangzhaohui@mail.tsinghua.edu.cn](mailto:wangzhaohui@mail.tsinghua.edu.cn))

(the product of molar extinction coefficient and  $\Phi_{\text{FL}}$ , which is a crucial factor for enhancing imaging resolution) remains a significant challenge owing to the rapid internal conversion (IC) quenching processes, commonly known as the energy gap law [25].

Rylene diimides, a well-known class of polycyclic aromatic hydrocarbons, have garnered ever-increasing attention owing to their unique rigidly conjugated framework and excellent physicochemical properties, enabling their wide application in organic optoelectronic fields [26–43]. Among the rylene diimides family, perylene diimides (PDI) are renowned for their outstanding molar extinction coefficient ( $\epsilon$ ) and high  $\Phi_{\text{FL}}$  in the visible-light region [44–46], making them highly sought-after in the field of bioimaging [47–50]. However, their counterparts, the longitudinal  $\pi$ -extended terrylene diimides (TDI) and quaterrylene diimides (QDI) fluorophores, which possess deep-red/near-infrared absorption/emission and higher  $\epsilon$  values than PDI [51,52], suffer from a notable decrease in  $\Phi_{\text{FL}}$  and fluorescent brightness. Although considerable progress have been made in synthetic and biological explorations concerning TDI and QDI fluorophores [53,54], achieving high deep-red/near-infrared  $\Phi_{\text{FL}}$  and fluorescent brightness in TDI and QDI fluorophores remains a great challenge [55–59]. Nevertheless, recent reports have presented a growing body of evidence on organic fluorophores containing heavy atoms that exhibit an intriguing anti-heavy-atom effect, where the presence of heavy atoms enhances their  $\Phi_{\text{FL}}$  and fluorescent brightness, contrary to conventional understanding [60–64]. Although the anti-heavy-atom effect of transition metals [61–66] and halogens [67–69] has been studied in a few instances in the literature, the potential anti-heavy-atom effect of chalcogens remains unexplored.

Herein, we successfully synthesized a series of TDI and QDI fluorophores through straightforward nitration and chalcogen-annulation in the bay region (Scheme 1). Our research revealed the chalcogen-annulation strategy can considerably enhance the deep-red/near-infrared  $\Phi_{\text{FL}}$  and fluorescent brightness of TDI and QDI fluorophores up to 10 times (Figure 1a). It is in stark contrast with the decrease in  $\Phi_{\text{FL}}$  and fluorescent brightness observed in the chalcogen-annulation of the PDI fluorophores (Figure 1b) [37]. To the best of our knowledge, this is the first example of the anti-heavy-atom effect observed with chalcogens ( $\Phi_{\text{FL}}(\text{O-TDI}) < \Phi_{\text{FL}}(\text{S-TDI}) < \Phi_{\text{FL}}(\text{Se-TDI})$ ) and the  $\Phi_{\text{FL}}$  (93%) and fluorescent brightness of Se-TDI ( $128,200 \text{ cm}^{-1} \text{ mol}^{-1} \text{ L}$ ) ranks among those of the top deep-red/near-infrared organic fluorophores available [70]. The time-resolved fluorescence lifetime and fs-TA experiment reveal that the fluorescence lifetimes are below 3.65 ns and the absence of obvious changes of the excited state lifetimes in TDI and QDI fluorophores. These results indicate that intersystem crossing (ISC) can be neglected in TDI and QDI fluorophores. The-



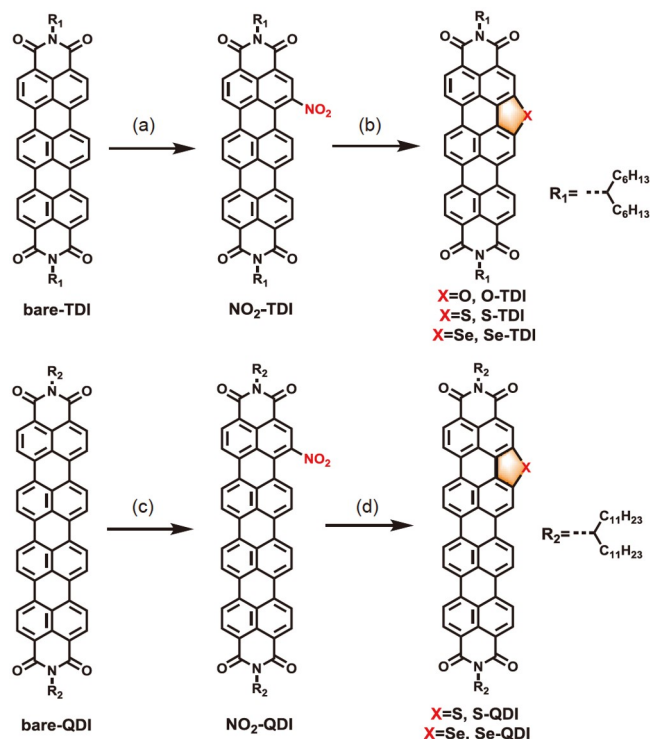
**Figure 1** (a) Change rules in  $\Phi_{\text{FL}}$  and fluorescent brightness of TDI and QDI fluorophores via chalcogen-annulation. (b) Heavy-atom and anti-heavy-atom effect in chalcogen-annulated PDI and TDI. (c) Jablonski diagram of energy conversion processes of TDI fluorophores (color on-line).

oretical calculations further confirm these results and reveal the suppressing of internal conversion (IC) processes of TDI fluorophores by reducing the reorganization energy of several important accepting modes in the ground state (Figure 1c). Our chalcogen-annulation strategy, which effectively increases the  $\Phi_{\text{FL}}$  and restricts the IC processes, while remaining unaffected by heavy-atom effect, offers novel insights and theoretical support for the design and synthesis of deep-red/near-infrared organic fluorophores with high  $\Phi_{\text{FL}}$  and fluorescent brightness.

## 2 Results and discussion

### 2.1 Synthesis of chalcogen-annulated TDI and QDI

The choice of nitration reagents plays a crucial role in the nitration reaction, as TDI and QDI demonstrate high susceptibility to nitration (Scheme 1). The use of nitrating reagents such as nitric acid can lead to the formation of dinitro- and trinitro-products within a short timeframe, diminishing the yield of mononitro-products and complicating the separation process [71]. Through extensive experimentation with various nitrating reagents, we ultimately opted for a relatively mild nitration reagent, tert-butyl nitrite, for the nitration reaction. By fine-tuning the reaction conditions, we achieved impressive yields of up to 90% for both  $\text{NO}_2$ -TDI



**Scheme 1** Synthesis routes for **O-TDI**, **S-TDI**, **Se-TDI**, **S-QDI** and **Se-QDI**. (a) Tert-butyl nitrite,  $\text{CH}_2\text{Cl}_2$ , r.t., 4 h, 90%. (b) **O-TDI**: air, NMP, 190 °C, 12 h, 15%; **S-TDI**: S powder, NMP, 190 °C, 2 h, 60%; **Se-TDI**: Se powder, NMP, 190 °C, 4 h, 50%. (c) Tert-butyl nitrite,  $\text{CHCl}_3$ , r.t., 4 h, 90%. (d) **S-QDI**: S powder, NMP, 190 °C, 4 h, 50%; **Se-QDI**: Se powder, NMP, 190 °C, 4 h, 50% (color online).

and **NO<sub>2</sub>-QDI** using  $\text{CH}_2\text{Cl}_2$  and  $\text{CHCl}_3$  as solvents, respectively. The conditions for chalcogen-annulation reactions of **NO<sub>2</sub>-TDI** closely resemble those of **NO<sub>2</sub>-PDI** [37]. By employing *N*-methyl-2-pyrrolidone (NMP) as the solvent, air, S powder, and Se powder as the raw materials, we conducted the reactions at 180 °C for 12 h, 2 h, and 4 h, respectively. The yields of **O-TDI**, **S-TDI** and **Se-TDI** were 15%, 60%, and 50%, respectively. Using identical reaction conditions, we also successfully synthesized **S-QDI** and **Se-QDI**, both yielding 50%. Regrettably, our efforts to synthesize **O-QDI** were unsuccessful. The challenge stems from a notable decrease in reactivity observed in the annulation reaction between oxygen and longitudinally  $\pi$ -extended rylene diimides. This limitation might also account for the low yield of **O-TDI** at 15%, notably lower than the previously reported 30% yield for **O-PDI** [37]. The structures of **O-TDI**, **S-TDI**, **Se-TDI**, **S-QDI** and **Se-QDI** were unambiguously characterized using NMR and high-resolution mass spectra. These compounds have good solubility and stability in common organic solvents such as toluene, dichloromethane, tetrahydrofuran, dioxane and NMP.

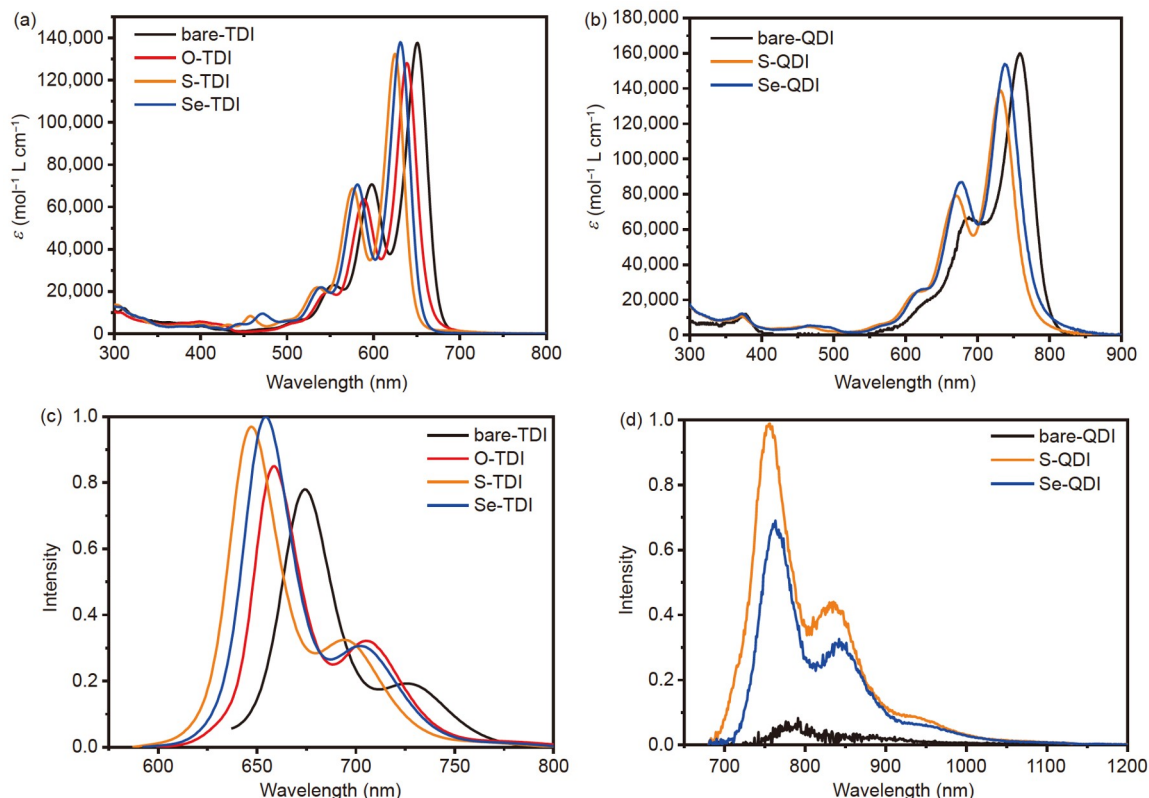
## 2.2 Absorption, fluorescence and electrochemical properties

The ground-state electronic properties of **TDI** fluorophores

were explored using steady-state absorption spectroscopy measurement in toluene at room temperature (Figure 2a–d). **Bare-TDI**, **O-TDI**, **S-TDI** and **Se-TDI** exhibit electronic absorption spectra corresponding to the  $S_0 \rightarrow S_1$  transition with well-resolved vibronic bands having absorption maxima at ( $\lambda_{\text{Abs max}}$ ) 650 nm for **bare-TDI**, 638 nm for **O-TDI**, 624 nm for **S-TDI**, and 631 nm for **Se-TDI**. The  $S_0 \rightarrow S_1$  electronic transition of **TDI** fluorophores is strongly coupled to the vinyl stretching mode of the **TDI** core, ensuring distinct vibronic progression in the absorption spectra (Figure 2a). The vibronic features are at 598 nm ( $\lambda_{\text{Abs 0-1}}$ ) for **bare-TDI**, 589 nm ( $\lambda_{\text{Abs 0-1}}$ ) for **O-TDI**, 576 nm ( $\lambda_{\text{Abs 0-1}}$ ) for **S-TDI**, and 582 nm ( $\lambda_{\text{Abs 0-1}}$ ) for **Se-TDI**. The hypsochromic shift of chalcogen-annulated **TDI** are 12, 26 and 19 nm with respect to **bare-TDI**. The variation pattern of the UV-vis-NIR absorption spectra of **QDI** fluorophores is very similar to that of **TDI** fluorophores, where the  $\lambda_{\text{Abs max}}$  of **bare-QDI**, **S-QDI** and **Se-QDI** is 790 nm, 766 nm and 775 nm (Figure 2b), respectively. And their  $S_0 \rightarrow S_1$  electronic transition is also strongly coupled to the vinyl stretching mode. And the vibronic features are at 685 nm ( $\lambda_{\text{Abs 0-1}}$ ) for **bare-QDI**, 670 nm ( $\lambda_{\text{Abs 0-1}}$ ) for **S-QDI**, and 676 nm ( $\lambda_{\text{Abs 0-1}}$ ) for **Se-QDI**. The hypsochromic shift of **S-QDI** and **Se-QDI** is 24 and 15 nm with respect to **bare-QDI**.

Steady-state fluorescence spectroscopy measurements were performed to understand the photoluminescence properties of chalcogenide **TDI** and **QDI** fluorophores in toluene. The fluorescence spectra of **bare-PDI**, **O-TDI**, **S-TDI**, and **Se-TDI** exhibit mirror-image characteristics with their respective steady-state absorption spectra (Figure 2c). The fluorescence emission peaks are observed at 667 nm for **bare-TDI** (Stokes shift,  $\Delta\nu'' \approx 392.1 \text{ cm}^{-1}$ ), 658 nm for **O-TDI** ( $\Delta\nu'' \approx 476.4 \text{ cm}^{-1}$ ), 638 nm for **S-TDI** ( $\Delta\nu'' \approx 351.7 \text{ cm}^{-1}$ ), and 645 nm for **Se-TDI** ( $\Delta\nu'' \approx 344.0 \text{ cm}^{-1}$ ). According to the Franck–Condon principle, we can identify the distinct vibronic features in the fluorescence spectra ( $\lambda_{\text{Em 0-1}}$ ) of **bare-TDI**, **O-TDI**, **S-TDI** and **Se-TDI** at 728 nm, 710 nm, 700 nm and 705 nm, respectively. These properties indicate that chalcogens-annulated **TDI** fluorophores are potential bioimaging dyes in deep-red regions. The fluorescence spectra of **QDI** fluorophores exhibit similar rigidity to **TDI** fluorophores, adhering to both the mirror-image rule and the Franck–Condon principle. The  $\lambda_{\text{Em max}}$  and  $\lambda_{\text{Em 0-1}}$  are at 790 nm and 860 nm for **bare-QDI**, 766 nm and 834 nm for **S-QDI**, 775 nm and 843 nm for **Se-QDI**, respectively (Figure 2d).

The absolute  $\Phi_{\text{FL}}$  of **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI** was determined to be 72%, 79%, 90%, and 93% (Table 1), respectively. Additionally, the steady-state absorption spectra reveal that their corresponding maximal  $\epsilon$  is 135,300, 127,700, 132,400 and 137,800  $\text{mol}^{-1} \text{ L cm}^{-1}$ , respectively. The progressive improvement of  $\Phi_{\text{FL}}$ , combined with the relatively constant maximal  $\epsilon$ , results in a gradual increase in



**Figure 2** UV-vis-NIR absorption and fluorescence emission spectra of **bare-TDI**, **O-TDI**, **S-TDI**, **Se-TDI**, **bare-QDI**, **S-QDI** and **Se-QDI** in  $10^{-5}$  mol L $^{-1}$  toluene solution. The intensity of fluorescence is determined using the ratio of absolute  $\Phi_{\text{FL}}$  (color online).

fluorescent brightness. The fluorescent brightness values of **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI** are 97,400, 100,900, 119,200, and 128,200 cm $^{-1}$  mol $^{-1}$  L, respectively. This consistent pattern of increasing  $\Phi_{\text{FL}}$  and fluorescent brightness signifies a notable anti-heavy atom effect. To our best knowledge, this represents the first reported instance of the anti-heavy atom effect of chalcogens. The chalcogen-annulated **QDI** also show improvements in  $\Phi_{\text{FL}}$  and fluorescent brightness. The  $\Phi_{\text{FL}}$  of **bare-QDI**, **S-QDI**, and **Se-QDI** was measured to be 1%, 10%, and 7% (Table 1), respectively. The corresponding maximal  $\varepsilon$  values of **bare-QDI**, **S-QDI**, and **Se-QDI** are 172,700, 150,100, and 155,400 mol $^{-1}$  L cm $^{-1}$ , respectively. As a result, the fluorescent brightness of **bare-QDI**, **S-QDI**, and **Se-QDI** gradually increases to 1,700, 15,000, and 10,900 mol $^{-1}$  L cm $^{-1}$ , respectively.

Cyclic voltammetry (CV) measurements were performed in CH $_2$ Cl $_2$  to investigate the electrochemical properties of the **TDI** and **QDI** fluorophores. The results, as shown in Table 1, Figures S1 and S2, reveal that both **TDI** and **QDI** fluorophores exhibit similar redox behavior characterized with one reversible oxidation wave and one reversible reduction wave. This indicates that the introduction of chalcogens to the **TDI** and **QDI** skeletons has negligible impact on their electrochemical properties. The half-wave reduction/oxidation potentials were  $-1.14/0.71$  V for **bare-TDI**,  $-1.17/0.78$  V for **O-TDI**,  $-1.19/0.76$  V for **S-TDI**,

$-1.20/0.75$  V for **Se-TDI**,  $-1.12/0.30$  V for **bare-QDI**,  $-1.18/0.36$  V for **S-QDI** and  $-1.17/0.39$  V for **Se-QDI** vs. Fc/Fc $^{+}$ . Therefore, the LUMO/HOMO/energy gap was determined to be  $-3.83/-5.36/1.53$  eV for **bare-TDI**,  $-3.82/-5.41/1.59$  eV for **O-TDI**,  $-3.82/-5.40/1.60$  eV for **S-TDI**,  $-3.81/-5.38/1.57$  eV for **Se-TDI**,  $-3.80/-5.00/1.20$  eV for **bare-QDI**,  $-3.75/-5.04/1.29$  eV for **S-QDI** and  $-3.74/-5.04/1.30$  eV for **Se-QDI**.

### 2.3 Femtosecond transient absorption analysis and time-resolved fluorescence lifetime measurement

The excited-state dynamics of the **TDI** and **QDI** fluorophores were probed by femtosecond transient absorption (fs-TA) measurements in a toluene solution (0.2–0.3 OD at 595 nm for **bare-TDI** and **O-TDI**, 575 nm for **S-TDI**, 580 nm for **Se-TDI**, 667 nm for **bare-QDI**, **S-QDI** and **Se-QDI**). The spectral and kinetic components of the fs-TA spectra were analyzed by singular value decomposition analysis, followed by global analysis using the Glotaran program package [72]. The fs-TA spectra of **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI** in the initial few picoseconds after excitation display a negative ground-state bleach (GSB) and stimulated emission (SE) ranging from 500 to 750 nm (Figure 3a–3d) along with positive excited-state absorption (ESA) between 750 and 1,300 nm (Figure S3), correspond-

**Table 1** Photophysical and electrochemical properties of **bare-TDI**, **O-TDI**, **S-TDI**, **Se-TDI**, **bare-QDI**, **S-QDI** and **Se-QDI**<sup>a)</sup>

| Compound <sup>a)</sup> | $\lambda_{\text{abs max}}$<br>(nm) | $\epsilon$<br>(mol <sup>-1</sup> L cm <sup>-1</sup> ) | $\lambda_{\text{em}}$ (nm) | $\Phi_{\text{FL}}$ <sup>a),b)</sup><br>(%) | $B_{\text{F}}$<br>(mol <sup>-1</sup> L cm <sup>-1</sup> ) | $E_{\text{LUMO}}$ <sup>c),f)</sup><br>(eV) | $E_{\text{HOMO}}$ <sup>c),f)</sup><br>(eV) | $E_{\text{g}}$ <sup>d)</sup><br>(eV) | $\tau_{\text{F}}$<br>(ns) | $k_{\text{r}}$ <sup>e)</sup><br>(10 <sup>7</sup> s <sup>-1</sup> ) | $k_{\text{nr}}$ <sup>e)</sup><br>(10 <sup>7</sup> s <sup>-1</sup> ) |
|------------------------|------------------------------------|-------------------------------------------------------|----------------------------|--------------------------------------------|-----------------------------------------------------------|--------------------------------------------|--------------------------------------------|--------------------------------------|---------------------------|--------------------------------------------------------------------|---------------------------------------------------------------------|
| Bare-TDI               | 650                                | 135,300                                               | 667, 728                   | 72                                         | 97,400                                                    | -3.83                                      | -5.36                                      | 1.53                                 | 3.52                      | 20.5                                                               | 7.95                                                                |
| O-TDI                  | 638                                | 127,700                                               | 658, 710                   | 79                                         | 100,900                                                   | -3.82                                      | -5.41                                      | 1.59                                 | 3.64                      | 21.7                                                               | 5.77                                                                |
| S-TDI                  | 624                                | 132,400                                               | 638, 700                   | 90                                         | 119,200                                                   | -3.82                                      | -5.42                                      | 1.60                                 | 3.62                      | 24.9                                                               | 2.76                                                                |
| Se-TDI                 | 631                                | 137,800                                               | 645, 705                   | 93                                         | 128,200                                                   | -3.81                                      | -5.38                                      | 1.57                                 | 3.60                      | 25.8                                                               | 1.94                                                                |
| Bare-QDI               | 758                                | 172,700                                               | 790, 860                   | 1                                          | 1,700                                                     | -3.80                                      | -5.00                                      | 1.20                                 | 0.45                      | 2.22                                                               | 220                                                                 |
| S-QDI                  | 731                                | 150,100                                               | 766, 834                   | 10                                         | 15,000                                                    | -3.75                                      | -5.04                                      | 1.29                                 | 1.27                      | 7.87                                                               | 70.9                                                                |
| Se-QDI                 | 737                                | 155,400                                               | 775, 843                   | 7                                          | 10,900                                                    | -3.74                                      | -5.04                                      | 1.30                                 | 1.02                      | 6.86                                                               | 91.2                                                                |

<sup>a)</sup> Measured in toluene (1.0×10<sup>-5</sup> mol L<sup>-1</sup>). <sup>b)</sup> Determined by absolute quantum yield method. <sup>c)</sup> Calculated by the onset of the first reduction peak according to  $E_{\text{LUMO}} = -(4.8 + E_{\text{onset}}^{\text{re}})$  eV,  $E_{\text{HOMO}} = (4.8 + E_{\text{onset}}^{\text{ox}})$  eV. <sup>d)</sup>  $E_{\text{g}} = E_{\text{LUMO}} - E_{\text{HOMO}}$ . <sup>e)</sup> The rates of radiative and nonradiative processes were determined by  $k_{\text{r}} = \Phi_{\text{FL}}/\tau_{\text{F}}$  and  $k_{\text{nr}} = (1 - \Phi_{\text{FL}})/\tau_{\text{F}}$ . <sup>f)</sup> Measured in dichloromethane.  $B_{\text{F}}$ =fluorescent brightness.

ing to the singlet excited-state ( $S_1 \rightarrow S_n$ ;  $n > 1$ ) of **bare-TDI**, **O-TDI**, **S-TDI** and **Se-TDI** [73]. The decay associated difference spectra (DADS) of each **bare-TDI**, **O-TDI**, **S-TDI** and **Se-TDI** were extracted (Figure 3e–3h). Three lifetimes ( $\tau_1$ ,  $\tau_2$ , and  $\tau_3$ ) were extrapolated for **bare-TDI**, **S-TDI** and **Se-TDI**. Two lifetimes ( $\tau_2$ , and  $\tau_3$ ) were extrapolated for **O-TDI**. Presumably,  $\tau_1$  (1.08, 5.48 and 6.65 ps for **bare-TDI**, **S-TDI** and **Se-TDI**, respectively) are assigned to the solvation or vibration-mediated relaxation to the lowest vibronic-state of the  $S_1$ ,  $\tau_2$  (452, 262, 370 and 481 ps for **bare-TDI**, **O-TDI**, **S-TDI** and **Se-TDI**, respectively) corresponds to solvation of hot  $S_1$  state as well as direct relaxation to the ground state from these higher vibrational  $S_1$  state, and  $\tau_3$  (3.59, 3.34, 3.33 and 3.65 ns for **bare-TDI**, **O-TDI**, **S-TDI** and **Se-TDI**, respectively) corresponds to the decay of the lowest vibrational  $S_1$  state to the ground-state. Meanwhile, the observed singlet lifetimes are comparable to the fluorescence lifetime of these molecules (3.53, 3.64, 3.62 and 3.60 ns for **bare-TDI**, **O-TDI**, **S-TDI** and **Se-TDI**, respectively, as shown in Table 1 and Figure S4).

The excited-state dynamics of **QDI** fluorophores are very similar to that of **TDI** fluorophores. The longest excited-state lifetimes of **bare-QDI**, **S-QDI** and **Se-QDI** are 0.28, 1.21 and 1.49 ns (Figure S5), respectively, and the excited-state lifetimes are also comparable to the fluorescence lifetimes (0.45, 1.27 and 1.02 ns for **bare-QDI**, **S-QDI** and **Se-QDI**, respectively, as shown in Table 1 and Figure S6). Based on the fs-TA and time-resolved fluorescence lifetime experiments, it can be concluded that ISC process is negligible in the **TDI** and **QDI** fluorophores, that is, the presence of heavy atoms does not accelerate the process of ISC.

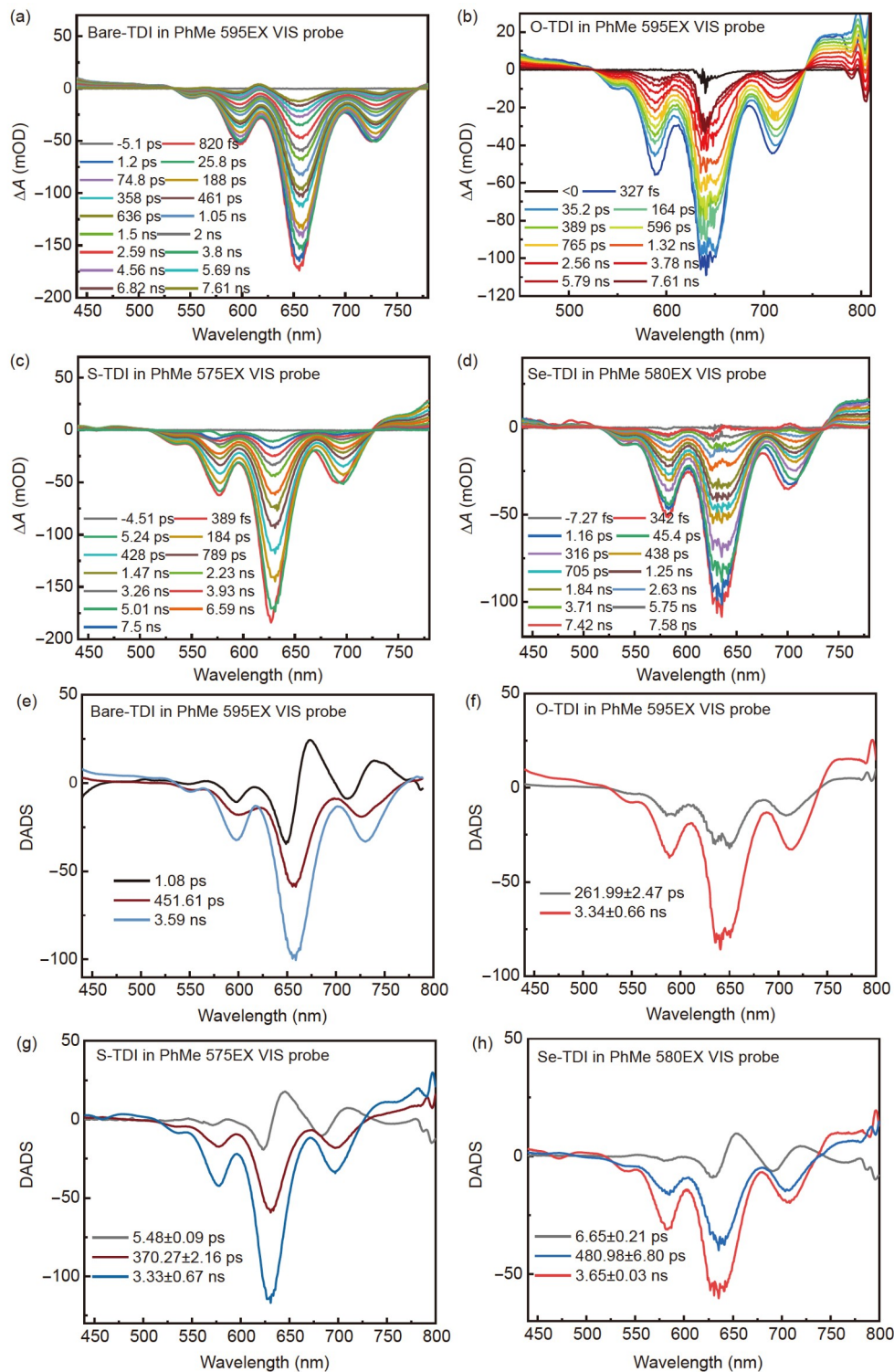
## 2.4 Theoretical calculations

To investigate the reason of enhancing the  $\Phi_{\text{FL}}$  and fluorescent brightness of **TDI** and **QDI** fluorophores, we conducted comprehensive theoretical calculations to thoroughly investigate the effect of annulation of chalcogens on these

processes. The rate of intersystem crossing ( $k_{\text{ISC}}$ ) is influenced by two crucial factors: the energy gap between the singlet- and triplet-states ( $\Delta E_{\text{ST}}$ ), and the total spin-orbit coupling matrix elements ( $H_{\text{SOC}}$ ) [74]. According to perturbation theory,  $k_{\text{ISC}}$  between singlet ( $S_m$ ) and triplet ( $T_n$ ) states is proportional to the square of their  $H_{\text{SOC}}$ . And according to the energy gap law,  $k_{\text{ISC}}$  is exponentially related to their  $\Delta E_{\text{ST}}$  (equation 1). The oscillator strength of singlet transitions confirms that the  $S_0 \rightarrow S_1$  transition is the most probable in **TDI** fluorophores (Table S1 and S2). Additionally, the  $T_3$  state is too high in energy to be reached in room temperature (Table S2 and S3). Therefore, the  $S_1 \rightarrow T_1$  and  $S_1 \rightarrow T_2$  transition may have the most significant impact on  $k_{\text{ISC}}$ . We employed the molecular materials property prediction package (MOMAP) [75–78] to calculate the rates of transitions from  $S_1$  to  $S_0$ ,  $T_1$  and  $T_2$ .

As shown in Table 2, the fluorescence rates of **bare-TDI** and chalcogen-annulated **TDI** are significantly larger than the internal conversion rates, leading to high  $\Phi_{\text{FL}}$ . The calculated rates of fluorescence ( $k_{\text{F}}$ , also called  $k_{\text{r}}$ ) exhibit an approximate 25% increase from **bare-TDI** to **O-TDI**, **S-TDI**, and **Se-TDI**, aligning with experimental trends. Considering the nearly identical oscillator strengths among the four molecules (Table S2) and minimal differences in  $S_1$ – $S_0$  energy gaps, it seems that changes in vibrational properties contribute to the increase in  $k_{\text{F}}$ . However, this observed increase remains marginal, and conducting a comprehensive analysis to ascertain the precise underlying causes of the enhanced  $\Phi_{\text{FL}}$  remains challenging.

On the other hand, the intersystem crossing rates are all negligible in **TDI** fluorophores. The large energy gap and negligible spin-orbital coupling matrix element between  $S_1$  and  $T_1$  make the  $S_1 \rightarrow T_1$  transition very slow. Meanwhile, the transition from  $S_1 \rightarrow T_2$  is also negligible, being very different from previously reported **PDI** systems [37]. The annulation of chalcogens in **TDI** fluorophores does not considerably alter the energy gap between  $S_1$  and  $T_2$  states, and the difference is higher than 0.25 eV in **bare-TDI**, **O-**



**Figure 3** Femtosecond transient absorption spectra of (a) **bare-TDI**, (b) **O-TDI**, (c) **S-TDI** and (d) **Se-TDI** in toluene. Reconstructed decay-associated difference spectra (DADS) of (e) **bare-TDI**, (f) **O-TDI**, (g) **S-TDI** and (h) **Se-TDI** extracted from global analysis (color online).

**TDI**, **S-TDI**, and **Se-TDI**. Generally, energy barriers that can be crossed at room temperature are usually below 0.2 eV. As a result, the transition from  $S_1$  to  $T_2$  in **TDI** fluorophores is nearly forbidden. Furthermore, the  $H_{\text{SOC}}$  of **TDI** fluorophores is less than  $0.1 \text{ cm}^{-1}$ . The high  $\Delta E_{\text{ST}}$  and very low

$H_{\text{SOC}}$  result in the inhibition of ISC (Table 2).

$$k_{\text{ISC}} \propto |\langle S_m | \hat{H}_{\text{SOC}} | T_n \rangle|^2 \exp(-\Delta E / kT) \quad (1)$$

To further analyze why  $\Delta E_{\text{ST}}$  and  $H_{\text{SOC}}$  of chalcogen-anulated **TDI** and **QDI** do not increase with the atomic

number of O, S, and Se, we conducted calculations of electron density changes in the excited-states of each molecule and generated charge density difference (CDD) data and plots (Figure 4, Figures S7 and S8, Table 2 and Table S1) [79–81]. The contribution ratio of CDD from O, S, and Se atoms in chalcogen-annulated **TDI** is provided in Table 2 and Figure 4. It is evident that the CDD of  $S_1$  and  $T_2$  states in the chalcogen-annulated **TDI** primarily distributes on the conjugated skeleton. The contribution of O, S, and Se atoms to the excited  $S_1$  (1.37, 2.56 and 2.81% for O, S and Se) and  $T_2$  state (0.92, 0.54 and 2.19% for O, S and Se) is minimal and there is no significant change in CDD near the O, S, and Se atoms. Chalcogen-annulated **QDI** exhibit similar CDD data as chalcogen-annulated **TDI** (Table S1 and Figure S7), the S and Se atomic orbitals do not play a significant role in the  $S_1 \rightarrow T_2$  transition process of chalcogen-annulated **QDI**. Consequently, we can conclude that the annulation of chalcogens does not significantly impact the properties of  $S_1$  and  $T_2$  states in chalcogen-annulated **TDI** and **QDI**, leading to negligible changes in  $\Delta E_{S_1-T_2}$  and  $H_{\text{SOC } S_1-T_2}$ .

It is noteworthy that we also calculated the CDD of chalcogen-annulated **PDI**, as depicted in Figure S8 and Table S1. We observe substantial contributions of S and Se to the CDD of chalcogen-annulated **PDI** at  $T_2$  state, reaching 31.5% and 42.5%, respectively. The high CDD of S and Se atoms leads to a significant reduction in the energy of the  $T_2$  state (reflected by the decrease of  $\Delta E_{S_1-T_2}$ ) and an increase in the

$H_{\text{SOC}}$  of  $S_1 \rightarrow T_2$  transition. Consequently, the  $k_{\text{ISC}}$  in chalcogen-annulated **PDI** is significantly enhanced. These findings are supported by previous reports where chalcogen-annulated **PDI** follows the heavy-atom effect [37]. The CDD calculations provide a plausible explanation for the distinct fluorescence properties of chalcogen-annulated **PDI** with chalcogen-annulated **TDI** and **QDI**.

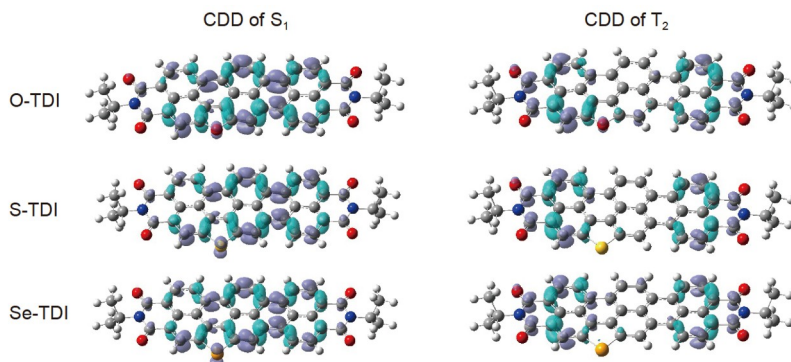
Interestingly, the internal conversion rates experience a substantial decrease by several folds. To investigate the impact of chalcogen-annulation on the IC process (the non-radiative transition from  $S_1$  to  $S_0$ ), we analyzed the ground-state reorganization energy of **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI**. The calculation results from MOMAP indicate negligible variation in the total reorganization energy of 769.0, 774.4, 790.3, and 783.2  $\text{cm}^{-1}$  for **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI**, respectively. However, the annulation of chalcogens enhances the rigidity of the annulative region and disrupts the molecular symmetry. This case leads to the dispersion of the large reorganization energy in **bare-TDI** into multiple smaller reorganization energies in **O-TDI**, **S-TDI**, and **Se-TDI**, leading to gradual decrease of the large reorganization energy in **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI** (255.56, 212.64, 157.58 and 148.44  $\text{cm}^{-1}$ , respectively), and as a result, gradually decreasing the rate of IC in **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI** (Table 2).

A more specific analysis is as shown in Figure 5a. In **bare-TDI**, the vibrational modes at 1,318  $\text{cm}^{-1}$  and 1,611  $\text{cm}^{-1}$

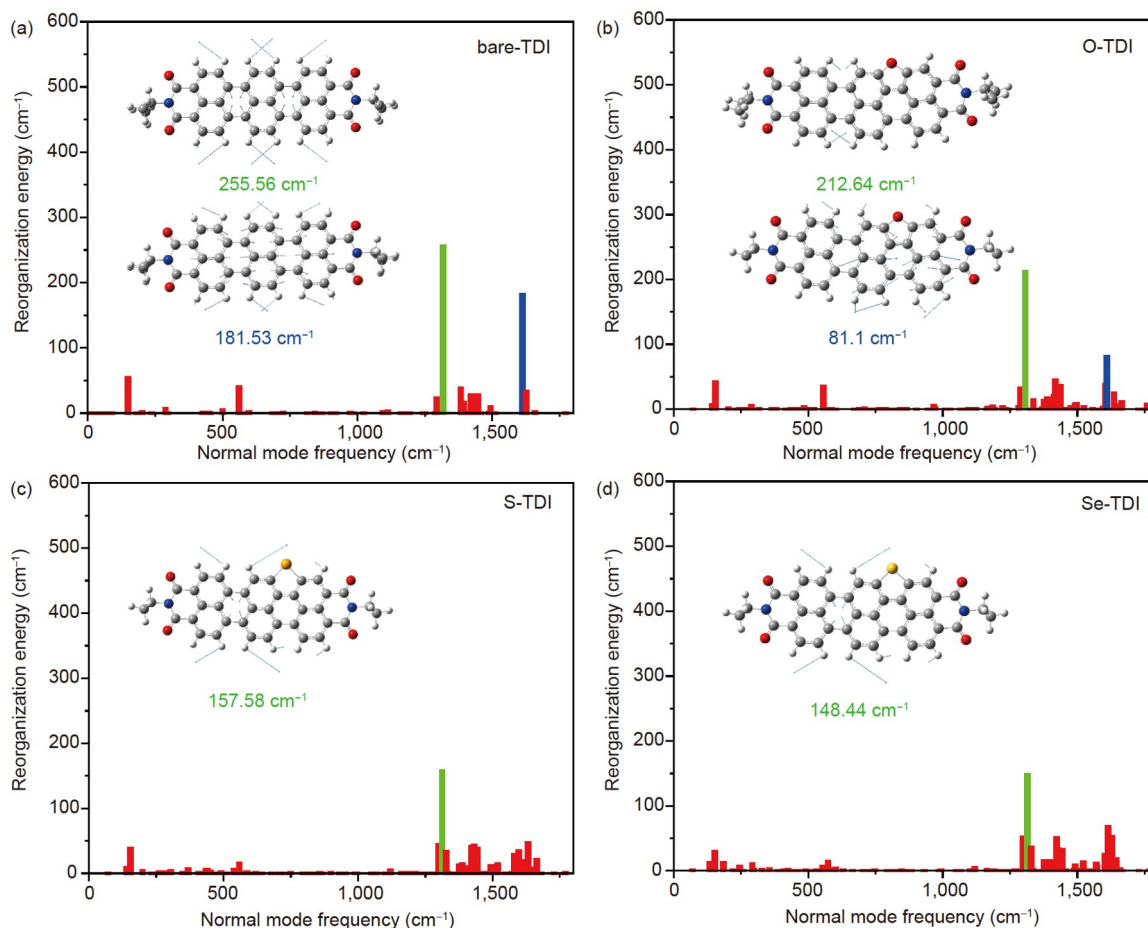
**Table 2** Calculated energies of  $S_1$ ,  $T_1$ ,  $T_2$ ,  $\Delta E_{\text{ST}}$ ,  $H_{\text{SOC}}$ ,  $k_{\text{ISC}}$  values of  $S_1 \rightarrow T_1$  and  $S_1 \rightarrow T_2$  transition, fluorescence rate ( $k_{\text{F}}$ ) and internal conversion rate ( $k_{\text{IC}}$ ) of  $S_1 \rightarrow S_0$  transition, in **Bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI**. The contribution of oxygen, sulfur and selenium atoms to CDD is also given. All energies in eV, except  $H_{\text{SOC}}$  in  $\text{cm}^{-1}$ . Rates in  $\text{s}^{-1}$

| Compound        | $S_1$ | $T_1$ | $T_2$ | $\Delta E_{\text{ST}}$<br>$S_1-T_2$ | $H_{\text{SOC}}$<br>$S_1-T_2$ | CDD<br>X% $S_1$ | CDD<br>X% $T_2$ | $k_{\text{F}}$<br>$S_1-S_0$ | $k_{\text{IC}}$<br>$S_1-S_0$ | $k_{\text{ISC}}$<br>$S_1-T_1^{\text{a}}$ | $k_{\text{ISC}}$<br>$S_1-T_2$ |
|-----------------|-------|-------|-------|-------------------------------------|-------------------------------|-----------------|-----------------|-----------------------------|------------------------------|------------------------------------------|-------------------------------|
| <b>Bare-TDI</b> | 1.61  | 0.52  | 1.86  | 0.25                                | 0.005                         | /               | /               | $1.30 \times 10^8$          | $4.10 \times 10^6$           | $9.13 \times 10^{-5}$                    | $1.53 \times 10^{-1}$         |
| <b>O-TDI</b>    | 1.69  | 0.65  | 1.95  | 0.26                                | 0.085                         | 1.37            | 0.92            | $1.29 \times 10^8$          | $8.00 \times 10^5$           | $3.08 \times 10^{-2}$                    | $3.06 \times 10^2$            |
| <b>S-TDI</b>    | 1.71  | 0.71  | 1.96  | 0.25                                | 0.087                         | 2.56            | 0.54            | $1.55 \times 10^8$          | $6.09 \times 10^5$           | $2.90 \times 10^{-3}$                    | $4.33 \times 10^2$            |
| <b>Se-TDI</b>   | 1.70  | 0.70  | 1.95  | 0.25                                | 0.092                         | 2.81            | 2.19            | $1.50 \times 10^8$          | $7.49 \times 10^5$           | 0                                        | $3.57 \times 10^2$            |

<sup>a)</sup> The Fermi golden rule may not be accurate for slow process such as the  $k_{\text{ISC}}$  of  $S_1-T_1$ , but it is certain that they are negligible.



**Figure 4** Charge density difference (CDD) plots of chalcogen-annulated **TDI** at  $S_1$  and  $T_2$  states at (TD)-B3LYP(GD3BJ)/def2svp level, the cyan represents the main distribution area of CDD (color online).



**Figure 5** Calculated reorganization energies of  $S_1 \rightarrow S_0$  are plotted against the normal-mode frequencies of (a) **bare-TDI**, (b) **O-TDI**, (c) **S-TDI** and (d) **Se-TDI**. High-intensity ( $>80 \text{ cm}^{-1}$ ) vibrational modes are shown as inset (color online).

have the highest reorganization energy, with values of  $255.56 \text{ cm}^{-1}$  and  $181.53 \text{ cm}^{-1}$ , respectively. They can be interpreted as longitudinal and transverse stretching vibrations of the two bridged benzene rings (Figure 5a inset). In **O-TDI**, the vibrational mode at  $1,318 \text{ cm}^{-1}$  changes to asymmetric longitudinal stretching vibrations primarily associated with the left hand-side benzene rings (upper inset of Figure 5b), while the right hand-side annulative region exhibits minimal vibration. Consequently, the reorganization energy of the vibrational mode at  $1,318 \text{ cm}^{-1}$  decreases from  $255.56 \text{ cm}^{-1}$  in **bare-TDI** to  $212.64 \text{ cm}^{-1}$  in **O-TDI**. The vibrational mode at  $1,611 \text{ cm}^{-1}$  splits into three vibrational modes of  $1,605.9$ ,  $1,609.9$  and  $1,664.3 \text{ cm}^{-1}$  in **O-TDI**, and the highest reorganization energy of the vibrational mode nearby  $1,611 \text{ cm}^{-1}$  also changes from  $181.53 \text{ cm}^{-1}$  in **bare-TDI** to  $81.1 \text{ cm}^{-1}$  in **O-TDI** (under inset of Figure 5b). This phenomenon becomes even more pronounced in **S-TDI** and **Se-TDI**, where the reorganization energy of these accepting modes at ground-state is  $157.58 \text{ cm}^{-1}$  and  $148.44 \text{ cm}^{-1}$ , respectively (Figure 5c and 5d). Overall, the chalcogen-annulation of **TDI** fluorophores with the progressive decrease in reorganization energy of several most important accepting

modes in the ground-state effectively suppresses the likelihood of IC from  $S_1$  to  $S_0$ . This inhibition becomes more pronounced with the increase in atomic number of the chalcogens (**O-TDI** < **S-TDI** < **Se-TDI**).

### 3 Conclusions

In summary, we demonstrated a straightforward and effective chalcogen-annulation strategy to enhance the deep-red/near-infrared  $\Phi_{\text{FL}}$  and fluorescent brightness of **QDI** and **TDI** fluorophores. The  $\Phi_{\text{FL}}$  of **S-QDI** and **Se-QDI** is 10%, and 7% and their fluorescent brightness values are 15,000, and  $10,900 \text{ cm}^{-1} \text{ mol}^{-1} \text{ L}$ . These values represent a tenfold enhancement compared with those of **bare-QDI** ( $\Phi_{\text{FL}} = 1\%$ , and fluorescent brightness =  $1,700 \text{ cm}^{-1} \text{ mol}^{-1} \text{ L}$ ). The  $\Phi_{\text{FL}}$  (72%, 79%, 90% and 93%) and the fluorescent brightness ( $97,400$ ,  $109,000$ ,  $119,200$ , and  $128,200 \text{ cm}^{-1} \text{ mol}^{-1} \text{ L}$ ) of **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI**, respectively, demonstrate a progressive increase. The continuous enhancement of  $\Phi_{\text{FL}}$  and fluorescent brightness of **TDI** fluorophores exhibits the first example of the anti-heavy-atom effect of

chalcogens and the  $\Phi_{\text{FL}}$  (93%) and fluorescent brightness ( $128,200 \text{ cm}^{-1} \text{ mol}^{-1} \text{ L}$ ) of **Se-TDI** is among the top of currently available deep-red/near-infrared fluorophores. The detailed mechanism is depicted in Figure 1c. The high energy gap and low  $H_{\text{SOC}}$  between the  $S_1 \rightarrow T_1$  and  $S_1 \rightarrow T_2$  indicate a negligible ISC rate in **TDI** fluorophores (the cyan squiggly arrow in Figure 1c). The chalcogen-annulation-induced alteration in vibrational properties potentially leads to a gradual increase in radiative rates within **TDI** fluorophores (green, orange, blue and red straight arrows in Figure 1c, respectively). Additionally, the decrease of the highest reorganization energy in the ground-states across **bare-TDI**, **O-TDI**, **S-TDI**, and **Se-TDI** (green, orange, blue and red squiggly arrows in Figure 1c, respectively) reduces the rate of IC in **TDI** fluorophores. Our chalcogen-annulation strategy provides valuable insights and theoretical supports for designing and synthesizing novel deep-red/near-infrared organic fluorophores with high  $\Phi_{\text{FL}}$  and brightness. The synthesis of chalcogen-annulated pentarylene- and hexarylene diimide fluorophores with high near-infrared-II fluorescence is ongoing in our laboratory.

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